

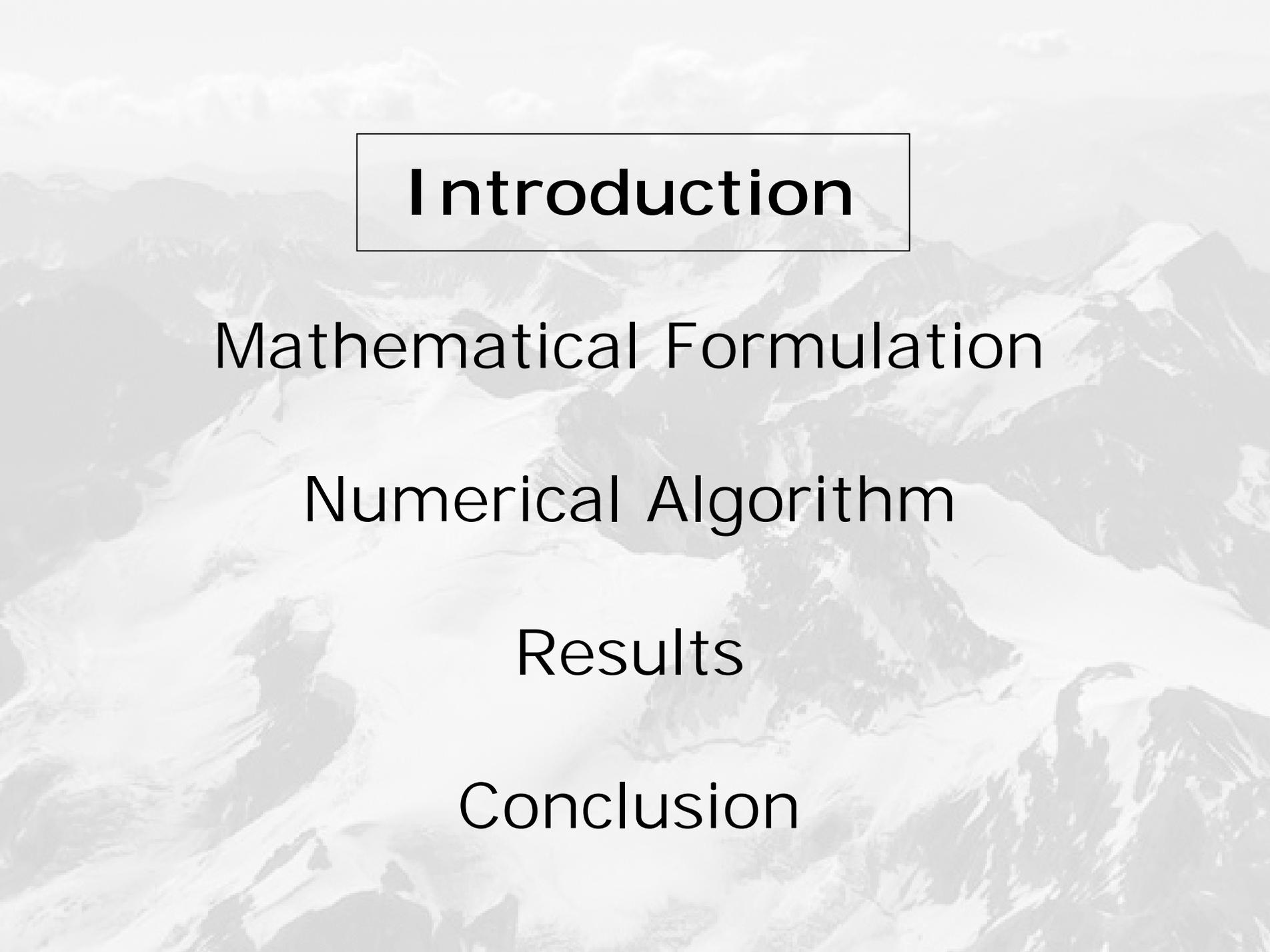
A Non-Hydrostatic Model to Simulate Atmospheric Flows in the Presence of Orography

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Outline of Presentation

- Introduction
- Mathematical Formulation
- Numerical Algorithm
- Results
- Conclusion



Introduction

Mathematical Formulation

Numerical Algorithm

Results

Conclusion

Motivation of the Research



Review of Previous Models

	Advection only	Shallow water equations	Compressible or hydrostatic models
No AMR		Qualitative model: <ul style="list-style-type: none">• wave equation	<ul style="list-style-type: none">• time step restricted by sound/gravity waves• terrain following coordinates• hydrostatic BVP is not well-posed
AMR	<ul style="list-style-type: none">• mass conservation equation only• no orography		Current research: <ul style="list-style-type: none">• time step restricted by advection only• EB formulation• well-posed BVP for AMR

AMR = Adaptive Mesh Refinement

Time steps

Typical cell: $\Delta x = 1.5\text{km}$ - $\Delta z = 200\text{m}$

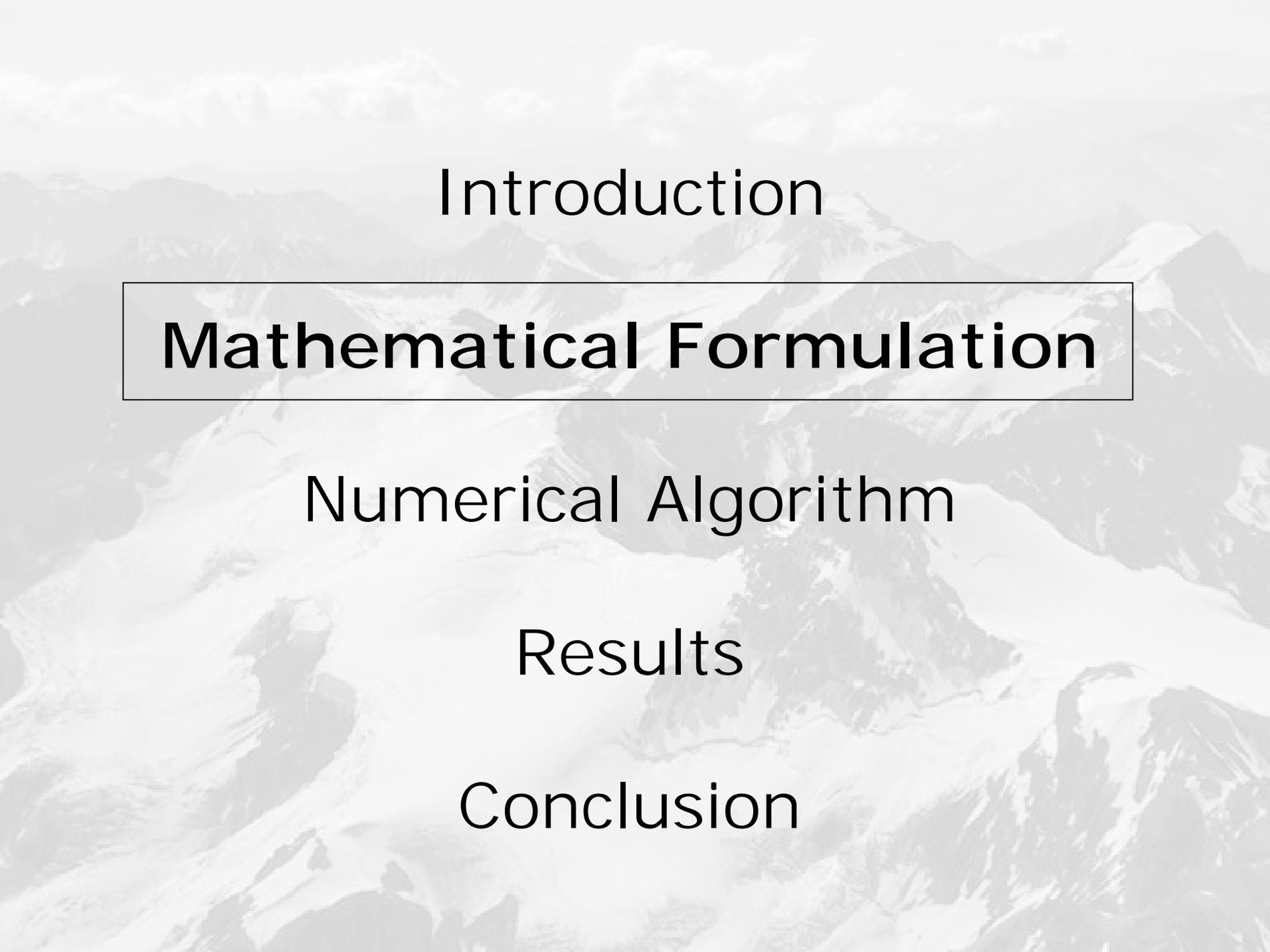
Time step limitation	Wave speed, c	$\Delta t = \frac{\{\Delta x, \Delta z\}}{c}$
Vertical acoustic waves	343 m/s	$\sim 0.6\text{ s}$
Horizontal gravity waves	200 m/s	$\sim 7.5\text{ s}$
Horizontal advection	20 m/s	$\sim 75\text{ s}$

Research Objective

- Develop a well-posed boundary value problem for gravitationally stratified flows to use in an AMR framework
- Applications
 - Atmospheric modeling
 - Astrophysics

Algorithmic Requirements

- Use advective time step
=> Implicit treatment of acoustic and gravity waves
- Adaptive Mesh Refinement (AMR)
=> Well-posed boundary-value formulation of the equations
- Orography
=> Cut-cell methods for irregular boundaries



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Euler Equations for a Compressible Fluid

Mass conservation

$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho \vec{u}) = 0$$

Momentum

$$\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \operatorname{grad}(\vec{u}) + \frac{1}{\rho} \operatorname{grad}(p) + g \vec{k} = \vec{0}$$

Pressure

$$\frac{\partial p}{\partial t} + \vec{u} \cdot \operatorname{grad}(p) + \rho c^2 \operatorname{div}(\vec{u}) = 0$$

Separating Out the Acoustic Waves: Hodge Decomposition

$$\mathbf{u} = \mathbf{u}_d + \mathbf{u}_p$$

total velocity incompressible vortical motions compressible motions

$\text{div}(\mathbf{u}_d) = 0$ $\mathbf{u}_p = \text{grad}(\phi)$

The diagram illustrates the Hodge decomposition of a vector field \mathbf{u} . It is shown as the sum of two components: \mathbf{u}_d (incompressible vortical motions) and \mathbf{u}_p (compressible motions). The total velocity \mathbf{u} is represented by a solid vertical line, while the decomposition is indicated by a dashed vertical line. The incompressible part \mathbf{u}_d is characterized by a divergence of zero, and the compressible part \mathbf{u}_p is represented as the gradient of a scalar potential ϕ .

Separating Out the Acoustic Waves: Projection Method

Incompressible flow and advective transport

$$\frac{\partial \tilde{\rho}}{\partial t} + \text{div}(\rho \vec{u}) = 0$$

$$\frac{\partial \vec{u}_d}{\partial t} + A_d \vec{u} + \frac{1}{\rho} \text{grad}(\pi) + P_o \left(\frac{1}{\rho} \text{grad} \delta + \frac{\tilde{\rho}}{\rho} g \vec{k} \right) = \vec{0}$$

Acoustic equations

$$\frac{\partial \vec{u}_p}{\partial t} + \text{grad} \left(\frac{|u_p + u_h|^2}{2} \right) + Q_o \left(\frac{1}{\rho} \text{grad} \delta + \frac{\tilde{\rho}}{\rho} g \vec{k} \right) = \vec{0}$$

$$\frac{\partial \delta}{\partial t} + \vec{u} \cdot \text{grad}(\pi + \delta) + \rho c^2 \text{div}(\vec{u}_p) + w \frac{\partial \rho_o}{\partial t} \vec{k} + \frac{\partial \pi}{\partial t} = 0$$

Incompressible flow => semi-implicit formulation

Acoustic waves => implicit formulation

Separating Out the Acoustic Waves: Projection Method

- Incompressible equations
 - Poisson-like equation for the pressure
 - Explicit treatment of advection
- Acoustic equations
 - Backward-Euler
 - Implicit treatment \Rightarrow Helmholtz equation for the acoustic pressure

\Rightarrow Well-posed boundary problems for AMR

Separating Out Fast Gravity Waves

- Isolate incompressible flow and gravity terms

$$\frac{\partial u_d}{\partial x} + \frac{\partial w_d}{\partial z} = 0$$

$$L_z \frac{\partial \pi_H}{\partial t} + \frac{\partial u_d}{\partial x} = f_\rho$$

$$\frac{\partial u_d}{\partial t} + \frac{1}{\rho_0} \frac{\partial \pi_H}{\partial x} = f_u$$

$$\frac{\partial \pi_H}{\partial z} = -\tilde{\rho}g$$

$\rho_0(z)$: background stratification

$L_z = \frac{\partial}{\partial z} \frac{1}{g} \frac{d\rho_0}{dz} \frac{\partial}{\partial z}$: second order self-adjoint operator

- This set of equations is equivalent to the set of incompressible equations in (u_d, ρ)

Separating Out Fast Gravity Waves

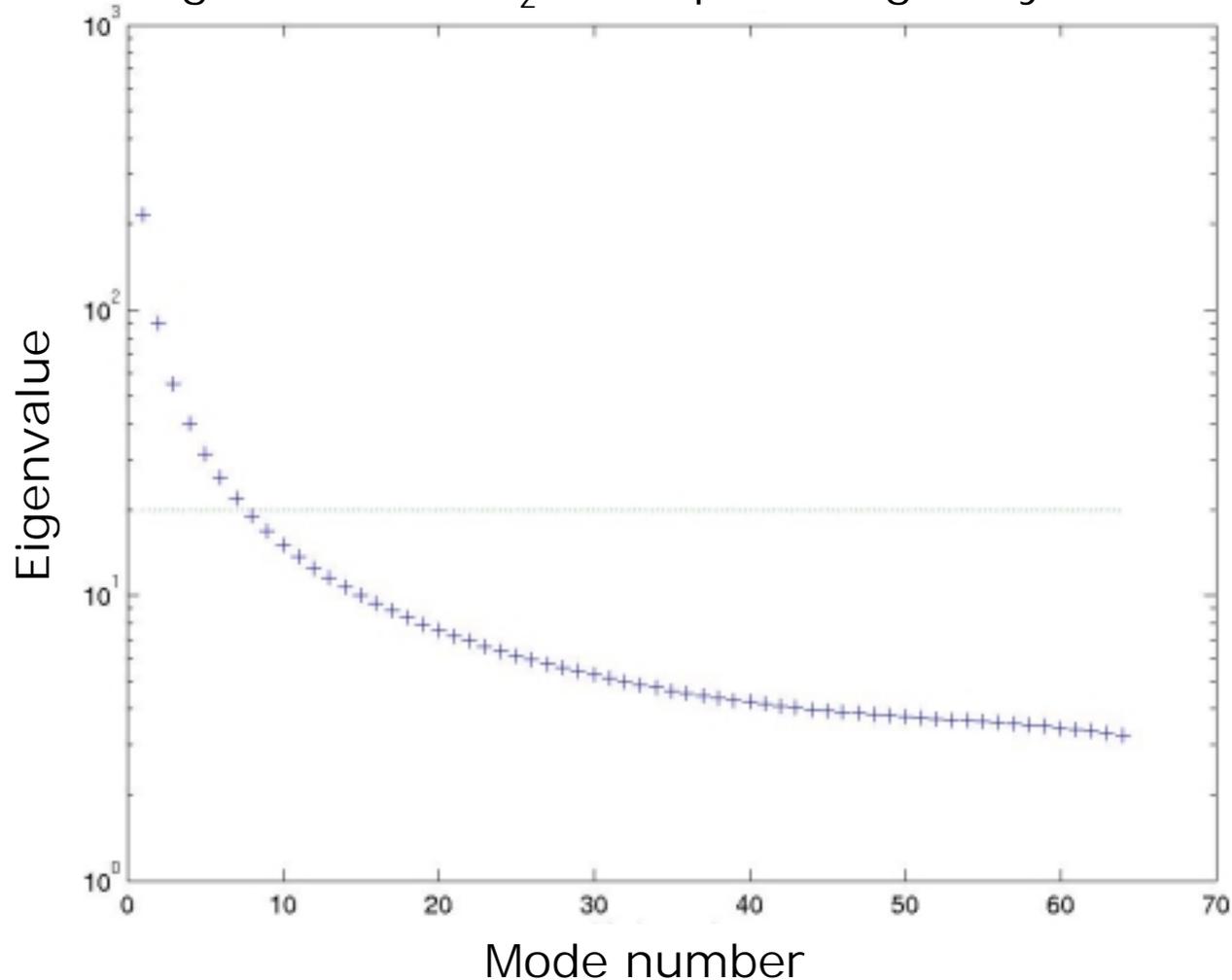
- Decomposition on eigenvectors

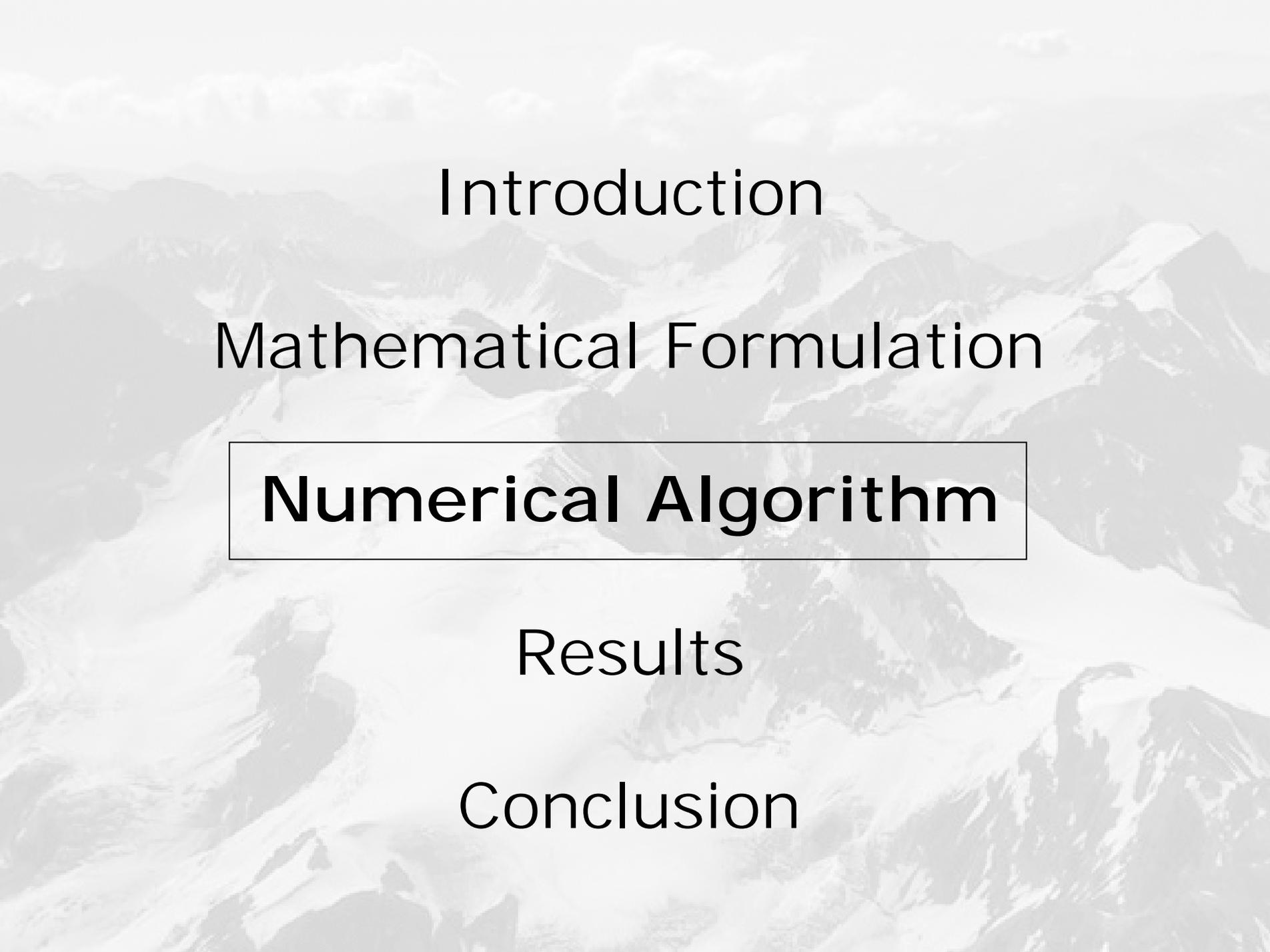
$$\begin{pmatrix} u_d \\ \pi_H \end{pmatrix} = \sum_k \begin{pmatrix} u_d^k(x, t) \\ \pi_H^k(x, t) \end{pmatrix} r^k(z)$$

- $\begin{pmatrix} u_d^k \\ \pi_H^k \end{pmatrix}$ satisfy the wave equation in x with wave speed $\lambda_k =$ eigenvalues of $L_z^{-1/2}$

Separating Out Fast Gravity Waves

Eigenvalues of $L_z^{-1/2}$ = speed of gravity waves





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Results

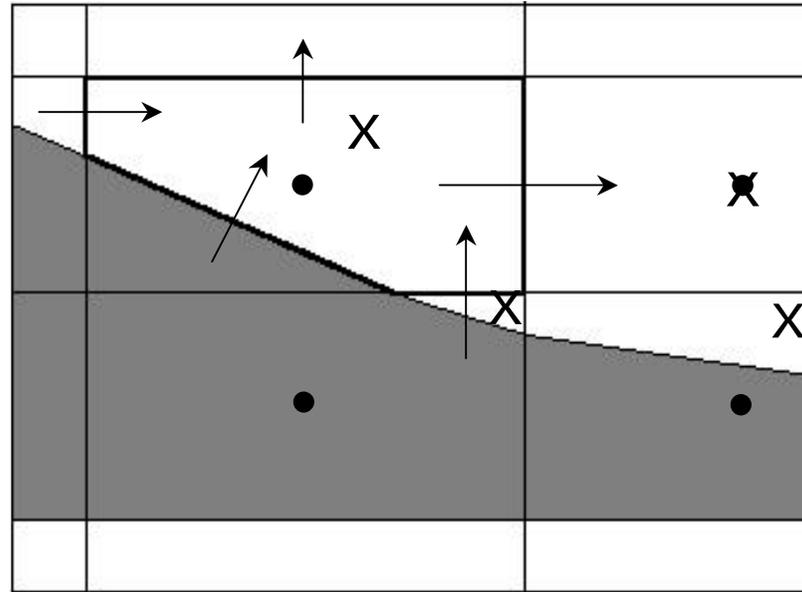
Conclusion

Overview of the Algorithm

$$\frac{\partial \vec{V}}{\partial t} + \vec{A} \left(\vec{V}, \frac{\partial \vec{V}}{\partial x}, \frac{\partial \vec{V}}{\partial z} \right) = \vec{0}$$

- Implicit treatment of acoustic waves using splitting
- Semi-implicit (explicit for advection) for incompressible advection
- Splitting of fast horizontal gravity waves from dynamics
- Use embedded boundaries for orography

Cartesian Grid Embedded Boundary methods



- PDEs written in conservation form

$$\frac{\partial U}{\partial t} + \nabla \cdot F(u) = 0 \quad \nabla \cdot \vec{F} \approx \frac{1}{V} \int_V \nabla \cdot \vec{F} dV = \frac{1}{V} \oint_S \vec{F} \cdot \vec{n} dS$$

- Away from boundaries: standard finite-difference discretization

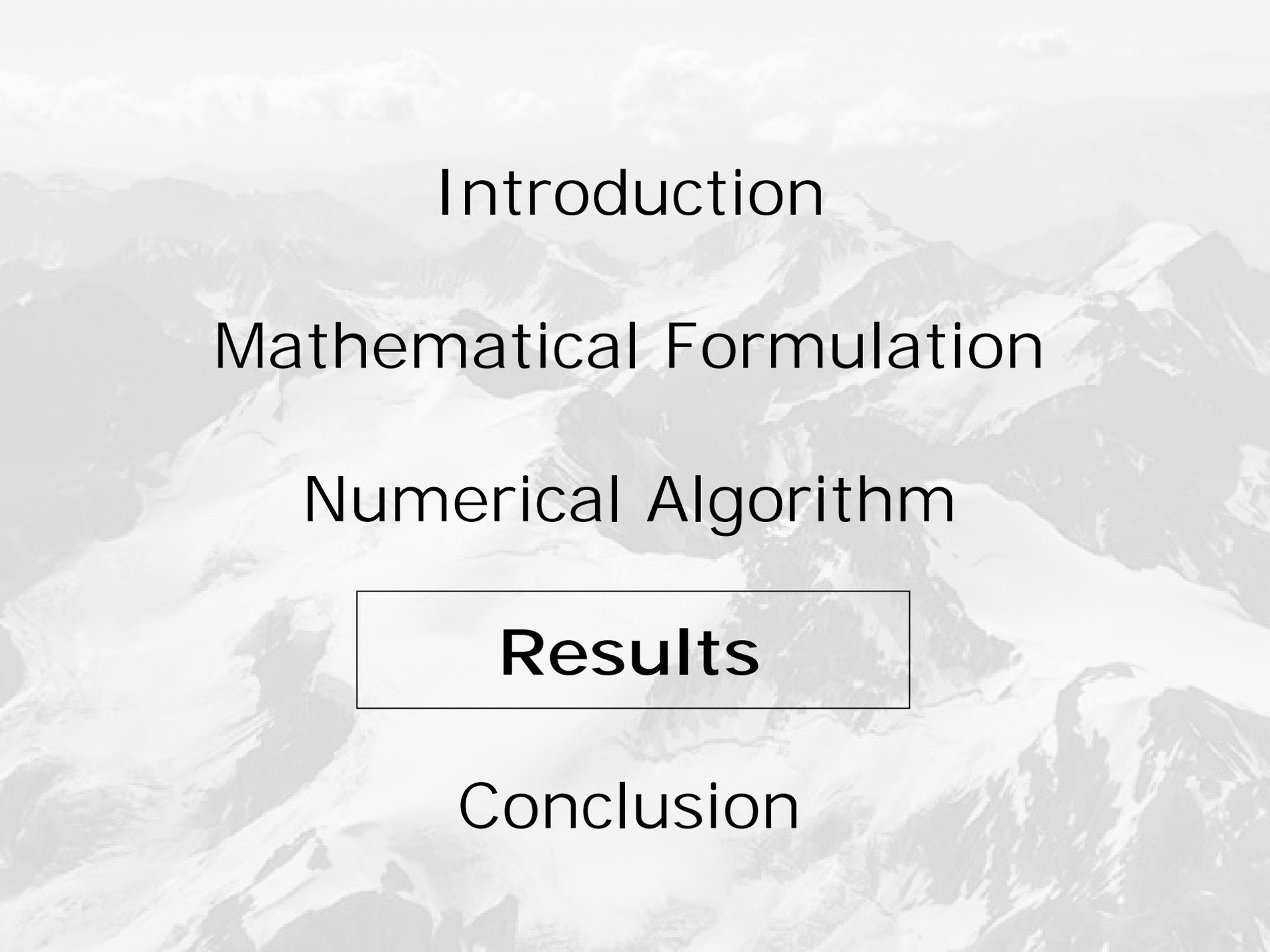
Cartesian Grid Embedded Boundary methods

Advantages of underlying rectangular grid

- Grid generation is tractable (T. Deschamps' talk)
- Well-understood
- Straightforward coupling to structured AMR

Large aspect ratio (1/10) introduces new issues

- Line solver for multigrid method



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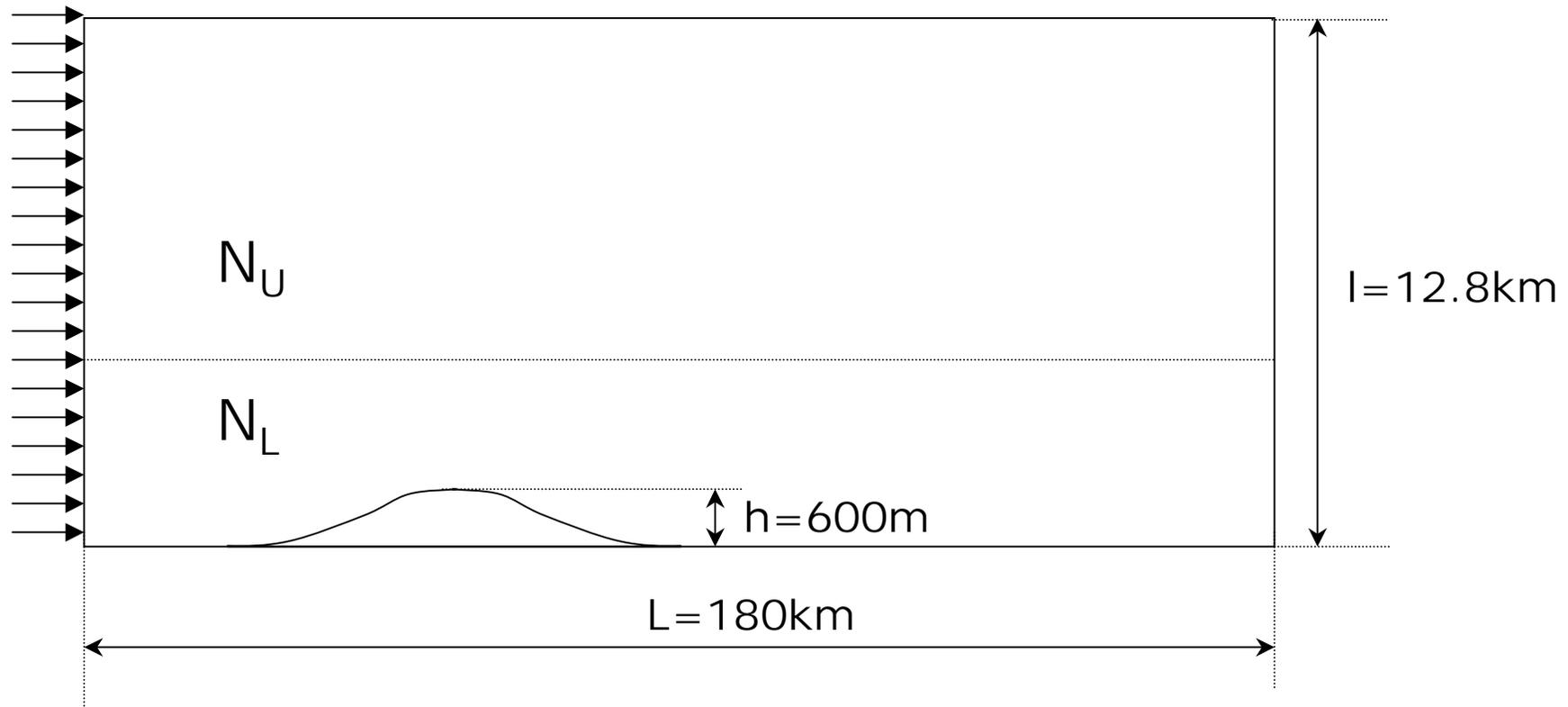
Results for a 2-Layered Atmosphere

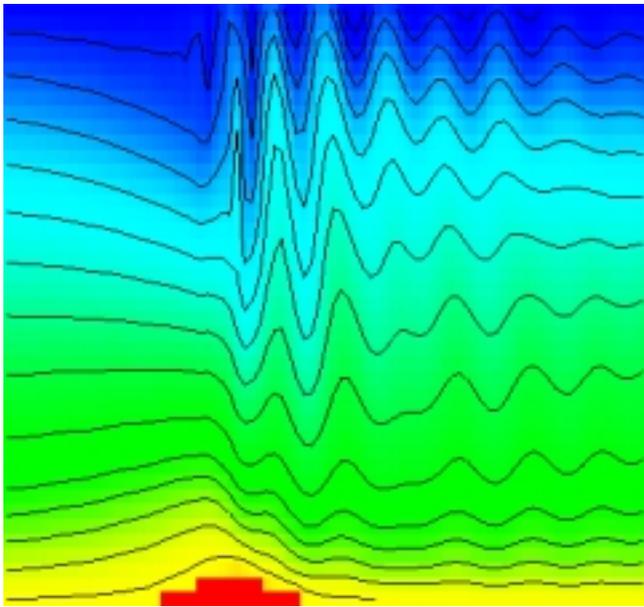
$\Delta x = 1406.25\text{m}$
 $\Delta z = 200\text{m}$
 $\Delta t = 5\text{s}$

N: Brunt-Väisälä frequency

- $N_U > N_L$: stable
- $N_L > N_U$: instable

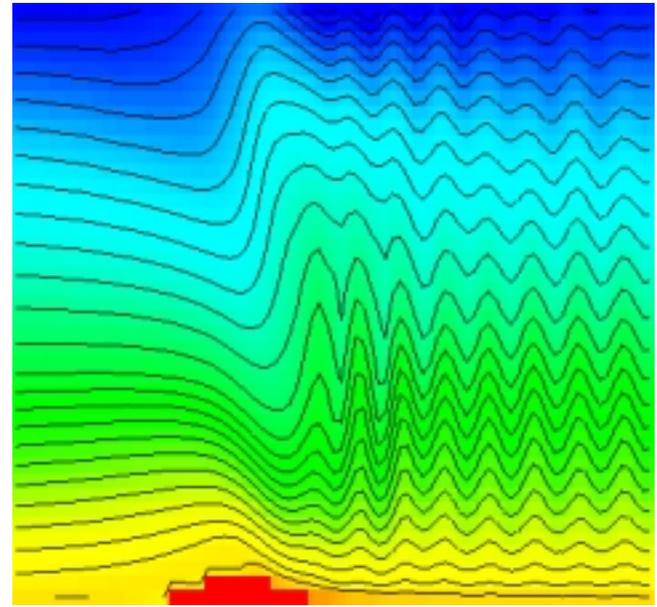
$u_0 = 20\text{m/s}$



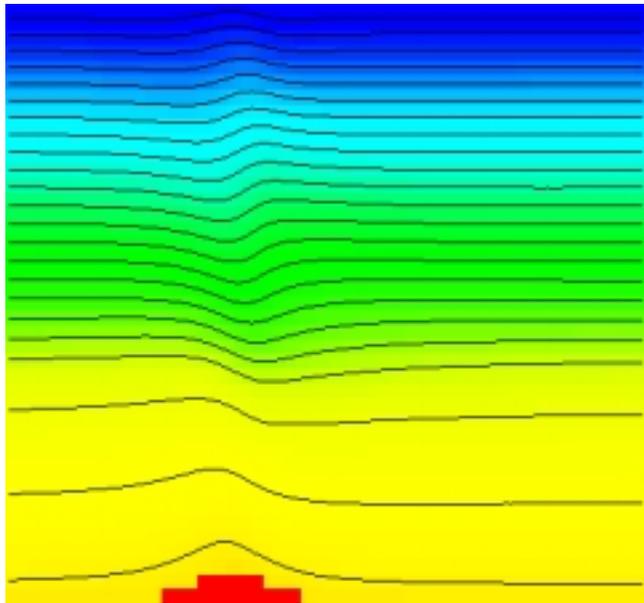


$$N_L > N_U$$

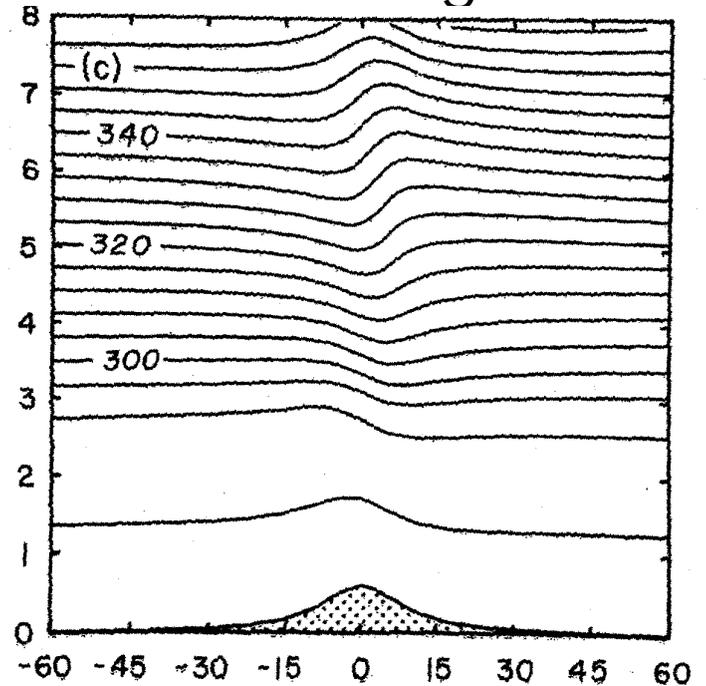
$\frac{1}{4}$ wavelength

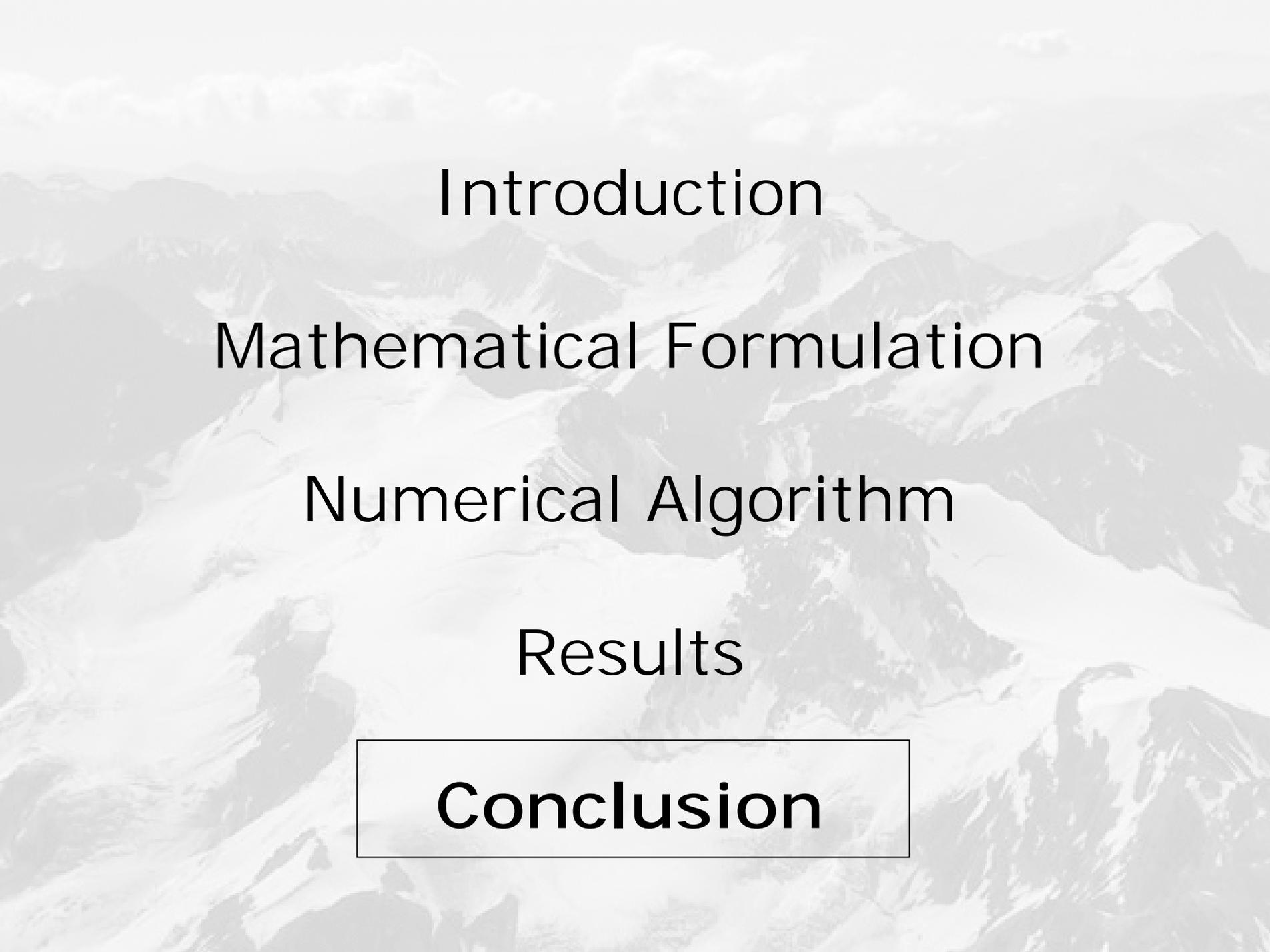


$\frac{1}{2}$ wavelength



$$N_U > N_L$$





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Status and Summary

- Formulation of well-posed initial-boundary problem
 - ✓ Stiff acoustic waves
 - ❑ Still working on implicit treatment of gravity waves
 - ❑ How does 2D implicit gravity waves couple with 3D AMR slow dynamics?
- EB treatment of orography
 - ✓ Good results on mountain lee-waves
- AMR
 - ❑ Not yet, but substantial progress on well-posed boundary-value problem

Future Work

- Split the gravity waves
- Go to 3D
- Use an AMR framework
- Parallelize the code



First-Aid Kit

Separating Out the Acoustic Waves: Hodge Decomposition

$$\begin{array}{c} 0 \\ \square \\ u_o \quad u \quad 0 \\ 0 \end{array} = \begin{array}{c} 0 \\ \square \\ 0 \quad u_d \quad 0 \\ 0 \end{array} + \begin{array}{c} 0 \\ \square \\ 0 \quad u_p \quad 0 \\ 0 \end{array} + \begin{array}{c} 0 \\ \square \\ u_o \quad u_h \quad 0 \\ 0 \end{array}$$

total velocity incompressible compressible

$\text{div}(u_d) = 0$ $\text{curl}(u_p) = 0$ $\text{div}(u_h) = 0$
 $\text{curl}(u_h) = 0$

Definitions of Projections

$$\vec{u}_d = P_o(\vec{u})$$

$$\vec{u}_\rho = Q_o(\vec{u})$$

$$Q_\rho = \frac{1}{\rho} \text{grad}(L_\rho)^{-1} \text{div}$$

$$P_\rho = I - Q_\rho$$

$$Q_o = \text{grad}(L_o)^{-1} \text{div}$$

$$P_o = I - Q_o$$

$$L_\rho = \text{div} \left(\frac{1}{\rho} \text{grad} \right)$$

$$L_o = \text{div}(\text{grad})$$

Definitions of Other Terms

$$A_d u = \bar{u} \cdot \text{grad}(\bar{u}) - \text{grad} \left(\frac{|u_p + u_h|^2}{2} \right)$$

$$\frac{1}{\rho} \text{grad}(\pi) = -Q_\rho(A_d u)$$

$$\delta = \rho - \rho_o(z) - \pi$$

Separating Out Fast Gravity Waves

$$\frac{\partial \pi_H}{\partial z} = -\bar{\rho}g$$

$$L_z \frac{\partial \pi_H}{\partial t} + \frac{\partial u_d}{\partial x} = f_\rho$$

$$\frac{\partial u_d}{\partial t} + \frac{1}{\rho_0} \frac{\partial \pi_H}{\partial x} = f_u$$

Wave equation

Project on fast eigenmodes

$$\lambda_k \frac{\partial \pi_H^k}{\partial t} + \frac{\partial \hat{u}_d^k}{\partial x} = \hat{f}_\rho$$

$$\frac{\partial \hat{u}_d^k}{\partial t} + \frac{\partial \pi_H^k}{\partial x} = \hat{f}_u$$

recompose

Back to original variables

$$\pi_H^{Fast} = \sum_{k=1}^K \pi_H^k r_k$$

$$\hat{u}_d^{Fast} = \sum_{k=1}^K \hat{u}_d^k r_k$$

$$u_d^{Fast}, \pi_H^{Fast}$$

Outline of Algorithm

Recall: unknowns are $u_d, u_p, \rho, \pi, \delta$

Find the eigenvalues and eigenvectors of L_z

1. Advance u_d to half time step and to face centers
2. Advance ρ to half time step and to face centers
3. Partially advance u_d using only the advective terms
4. Solve for π_H^{Fast}, u_d^{Fast}
5. Project out unstable (fast) modes
6. Solve for auxiliary pressure π
7. Advance acoustic pressure δ implicitly
8. Advance curl-free velocity u_p
9. Update ρ
10. Add missing terms in update for u_d

